



A NEW TREATMENT OF SINGULAR INTEGRALS IN GALERKIN BOUNDARY INTEGRAL EQUATIONS

MARC LENOIR & NICOLAS SALLES

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Outline

Introduction

The Method

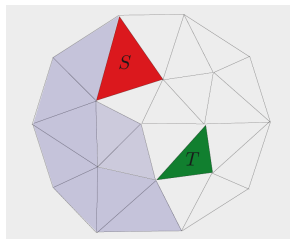
Comparison & Result

What do we want ?

We want to evaluate

$$\int_{\Gamma \times \Gamma} q(y) G(x, y) \bar{t}(x) d\gamma_x d\gamma_y \text{ and } \int_{\Gamma \times \Gamma} p(y) \frac{\partial G(x, y)}{\partial n_y} \bar{s}(x) d\gamma_x d\gamma_y,$$

We discretize Γ with plane polygons (usually triangles)



$$\int_{S \times T} \phi_i(x) G(x, y) \phi_j(y) d\gamma_x d\gamma_y$$

$$\int_{S \times T} \psi_i(x) \frac{\partial G(x, y)}{\partial n_y} \psi_j(y) d\gamma_x d\gamma_y$$

Singular Integrals

Green kernel and its gradient

$$G(x, y) = -\frac{e^{ik\|x-y\|}}{4\pi\|x-y\|} = -\frac{1}{4\pi} \frac{1}{\|x-y\|} + H(\|x-y\|)$$

$$\nabla G(x, y) = -\frac{1}{4\pi} \frac{x-y}{\|x-y\|^3} - \frac{k^2}{8\pi} \frac{x-y}{\|x-y\|} + (x-y)K(\|x-y\|)$$

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Singular integrals

$$\int_{S \times T} \frac{1}{\|x-y\|} dx dy; \int_{S \times T} \frac{x-y}{\|x-y\|^3} dx dy; \int_{S \times T} \frac{x-y}{\|x-y\|} dx dy$$

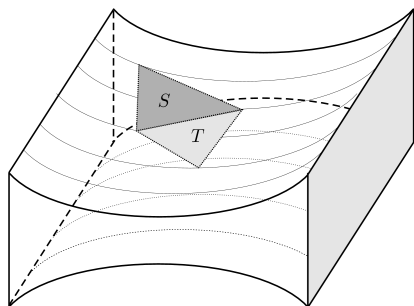
The issues

Everyone knows that it is essential to calculate these integrals very carefully and with high accuracy.

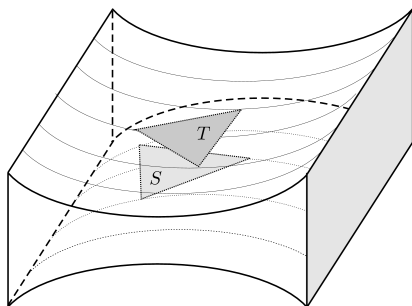
References

- M. G. Duffy. *Quadrature over a pyramid or cube of integrands with a singularity at a vertex*, SIAM J. Numer. Anal., 1982
- S. Nintcheu Fata, *Semi-analytic treatment of nearly-singular Galerkin surface integrals*, Journal of computational and applied mathematics, 2009
- S. Sauter and C. Schwab, *Boundary Element Methods*, Springer Series in Computational Mathematics, 2010

Singularities



Singular case



Almost singular case



Lenoir M., *Influence coefficients for variational integral equations*,
Comptes Rendus Mathematique, 2006

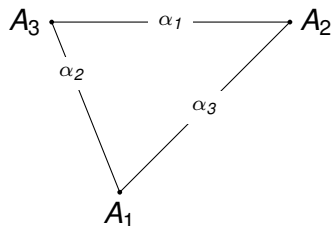
First Case : The self-influence

We want to evaluate **exactly** :

$$I = \int_{S \times S} \frac{dx dy}{\|x - y\|}$$

We define,

$$\sigma_i = \frac{\|P_i - A_i^-\|}{|\alpha_i|}$$



We will prove that we obtain

$$I = \frac{2|S|}{3} \sum_{i=1}^3 \gamma_i \left(\arg \sinh \left(\frac{|\alpha_i| - \sigma_i}{\gamma_i} \right) + \arg \sinh \left(\frac{\sigma_i}{\gamma_i} \right) \right).$$

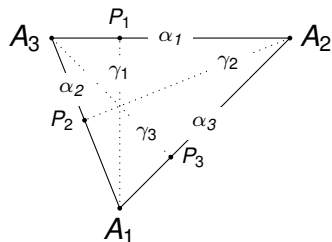
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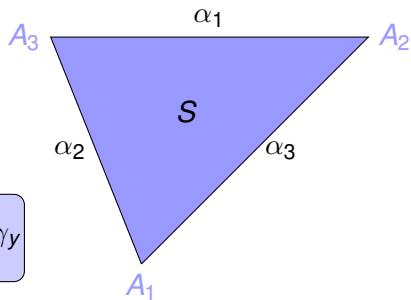
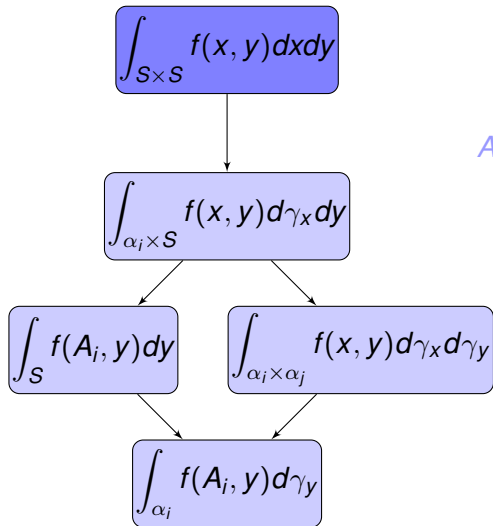
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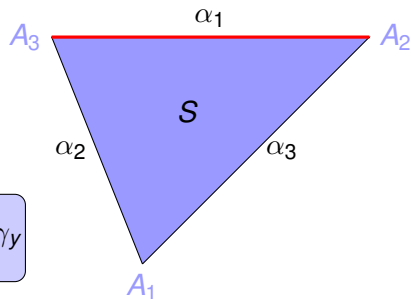
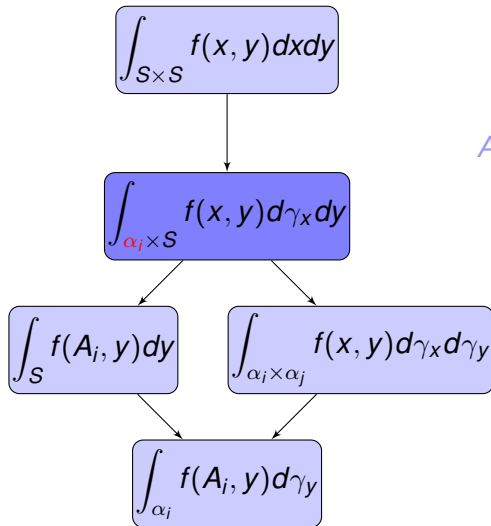
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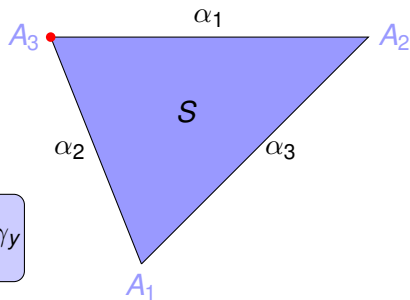
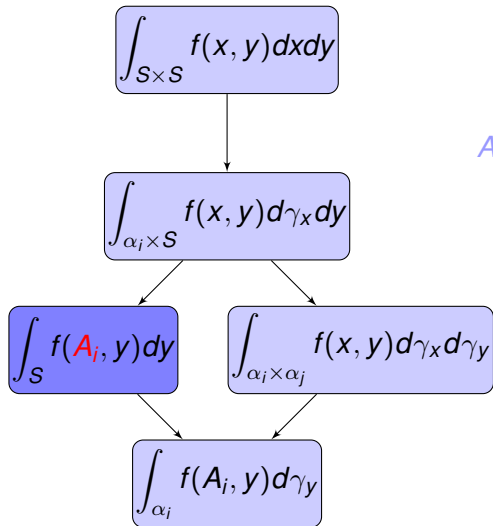
The Reduction Process



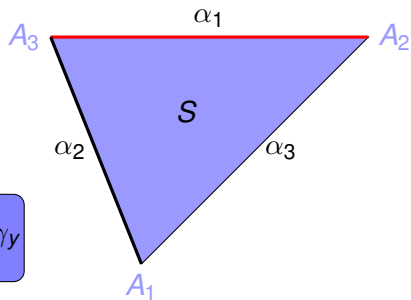
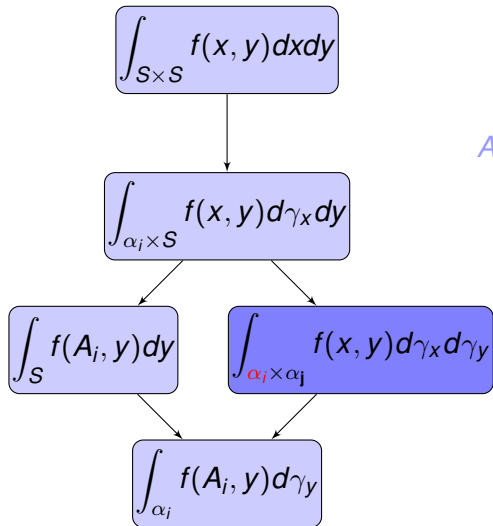
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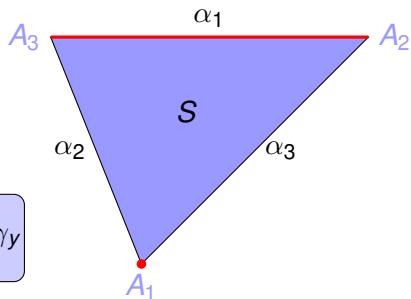
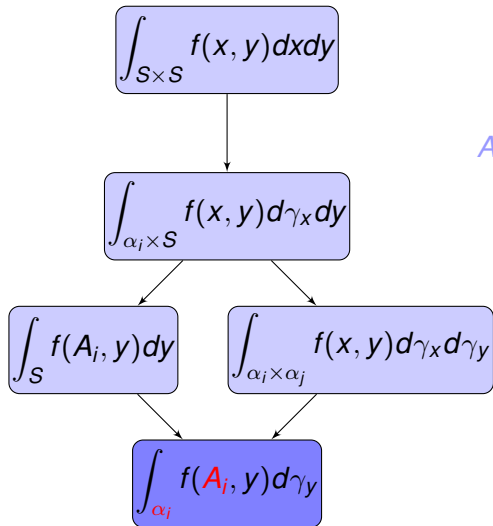
The Reduction Process



The Reduction Process



The Reduction Process



The First Formula

Suppose

- ▶ $\Omega \subset \mathbb{R}^n$
- ▶ $f : \Omega \rightarrow \mathbb{R}$ is a homogeneous function of degree r :

$$f(\lambda z) = \lambda^r f(z), \forall \lambda > 0$$

Then

$$\int_{\Omega} f(z) dz = \frac{1}{r+n} \int_{\partial\Omega} (z|\vec{\nu}) f(z) d\gamma_z$$

Proof.

By integration of the Euler's theorem on homogeneous functions □



Rosen D., Cormack D. E., *Singular and Near Singular Integrals in the Bem: A Global Approach*, SIAM J.A.M., 1993

The First Formula

In our case

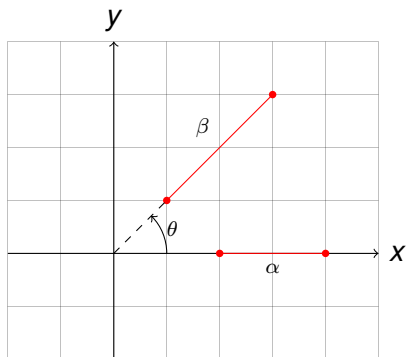
We use

$$\int_{\Omega} f(z) dz = \frac{1}{r+n} \int_{\partial\Omega} (z|\vec{\nu}) f(z) d\gamma_z$$

With

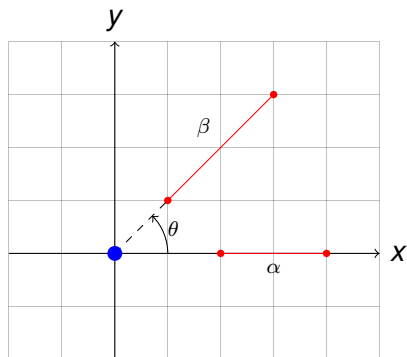
- ▶ $\Omega = \mathcal{S} \times \mathcal{S} \implies n = 4$
- ▶ $z = (x, y) \in \mathcal{S} \times \mathcal{S}$
- ▶ $f(x, y) = \frac{1}{\|x - y\|} \implies r = -1$
- ▶ Ω is a polyhedron so $\vec{\nu}$ is piecewise constant.

The choice of the origin



$$\int_{\alpha \times \beta} \frac{d\gamma_x d\gamma_y}{\|x - y\|}$$

The choice of the origin



$$x_1(s) = s$$

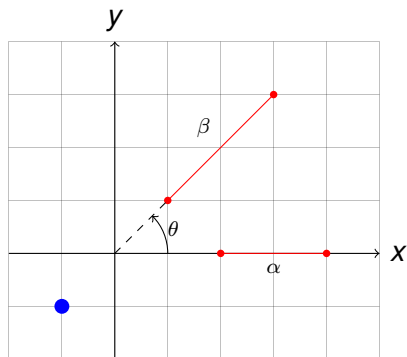
$$x_2(s) = 0$$

$$y_1(t) = t \cos(\theta)$$

$$y_2(t) = t \sin(\theta)$$

$$\int_{\alpha \times \beta} \frac{d\gamma_x d\gamma_y}{\|x - y\|} = \int_s \int_t \frac{ds dt}{\sqrt{(s - t \cos(\theta))^2 + (t \sin(\theta))^2}}$$

The choice of the origin



$$x_1(s) = s$$

$$x_2(s) = \mathbf{1}$$

$$y_1(t) = t \cos(\theta)$$

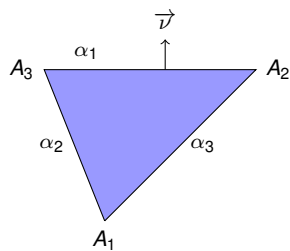
$$y_2(t) = t \sin(\theta)$$

$$\int_{\alpha \times \beta} \frac{d\gamma_x d\gamma_y}{\|x - y\|} = \int_s \int_t \frac{ds dt}{\sqrt{(s - t \cos(\theta))^2 + (\mathbf{1} - t \sin(\theta))^2}}$$

The reduction process

The self-influence

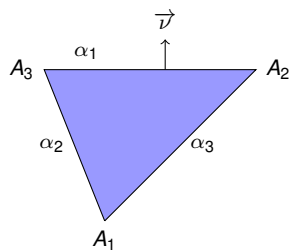
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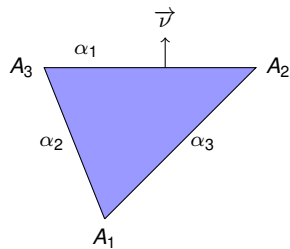


$$\int_{\Omega} f(z) dz = \frac{1}{r+n} \int_{\partial\Omega} (z|\vec{v}) f(z) d\gamma_z$$

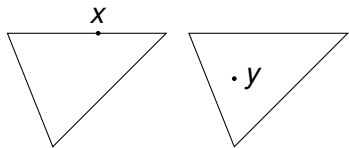
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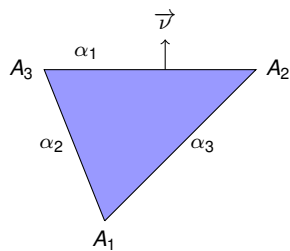
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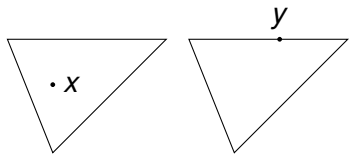
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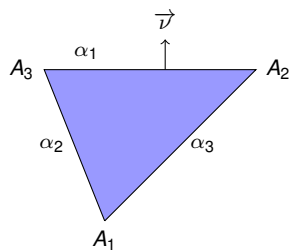
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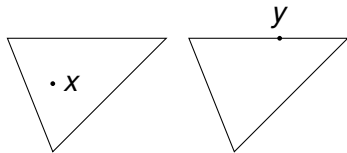
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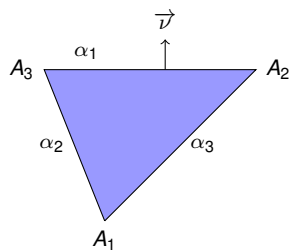
$$I = \frac{1}{3} \int_{\partial S \times S} \frac{(x | \vec{n})}{\|x - y\|} d\gamma_x dy + \frac{1}{3} \int_{S \times \partial S} \frac{(y | \vec{n})}{\|x - y\|} dx d\gamma_y$$



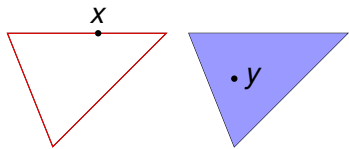
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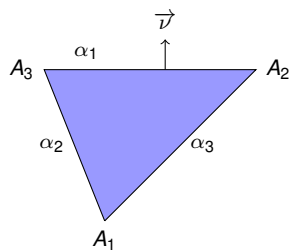
$$I = \frac{2}{3} \int_{\partial S \times S} \frac{(x | \vec{n})}{\|x - y\|} dxdy$$



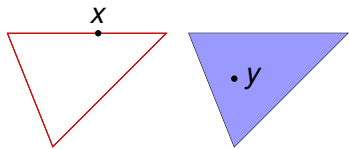
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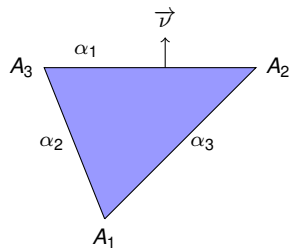
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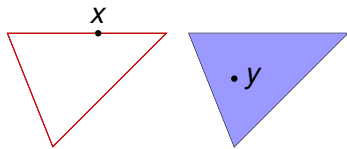
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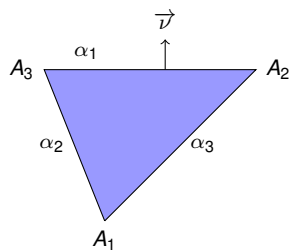
$$\partial S = \bigcup_{i=1..3} \alpha_i$$



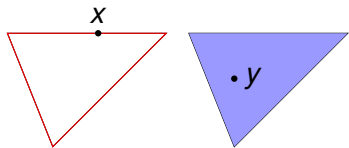
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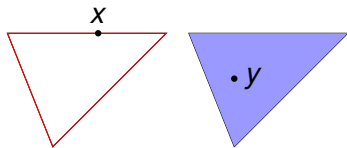
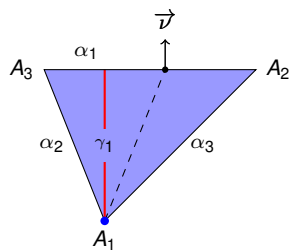
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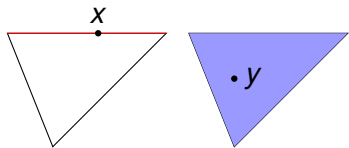
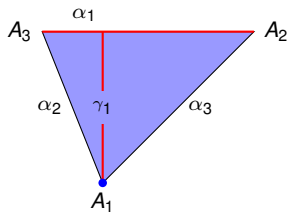
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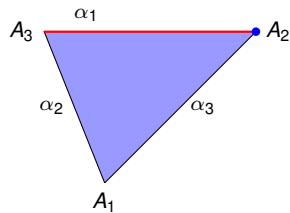
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The reduction process

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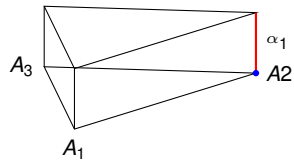
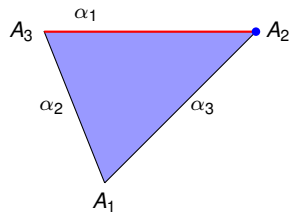
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The reduction process

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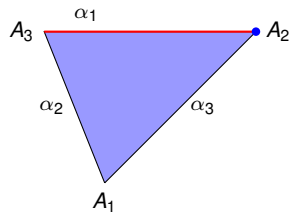
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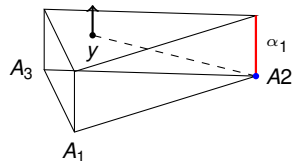
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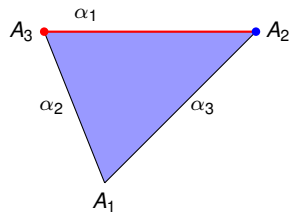
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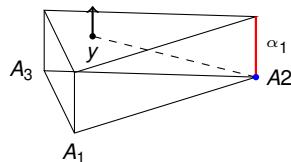
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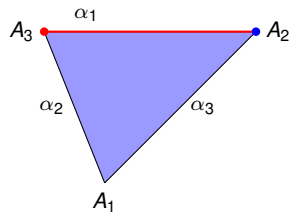
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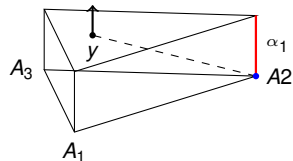
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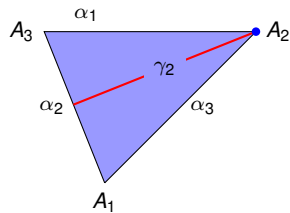
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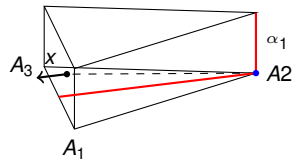
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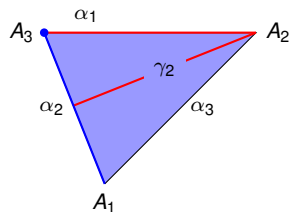
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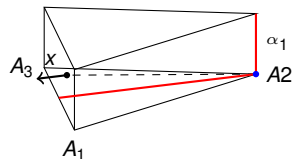
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$$I = \frac{2}{3} \gamma_1 \int_{\alpha_1 \times S} \frac{dxdy}{\|x - y\|}$$
$$I = \frac{2}{3} \gamma_1 \frac{|\alpha_1|}{2} \int_S \frac{dy}{\|A_3 - y\|}$$
$$+ \frac{2}{3} \gamma_1 \frac{\gamma_2}{2} \int_{\alpha_1 \times \alpha_2} \frac{d\gamma_x d\gamma_y}{\|x - y\|}$$



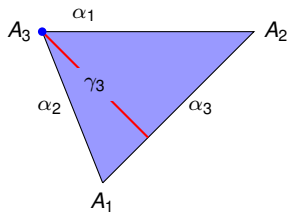
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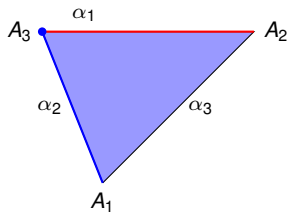
$$I = \frac{2}{3} |S| \int_S \frac{dy}{\|A_3 - y\|} + \frac{1}{3} \gamma_1 \gamma_2 \int_{\alpha_1 \times \alpha_2} \frac{d\gamma_x d\gamma_y}{\|x - y\|}$$



The reduction process

The self-influence

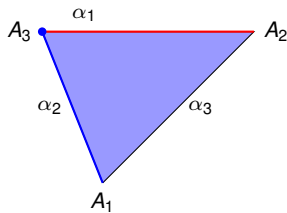
$$I = \frac{2}{3} |S| \gamma_3 \int_{\alpha_3} \frac{d\gamma_y}{\|A_3 - y\|} + \frac{1}{3} \gamma_1 \gamma_2 \int_{\alpha_1 \times \alpha_2} \frac{d\gamma_x d\gamma_y}{\|x - y\|}$$



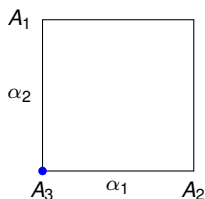
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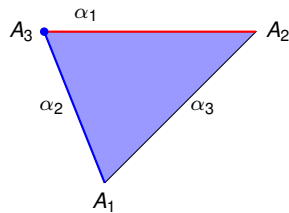
$$\partial(\alpha_1 \times \alpha_2) = (\{A_3, A_2\} \times \alpha_2) \cup (\alpha_1 \times \{A_3, A_1\})$$



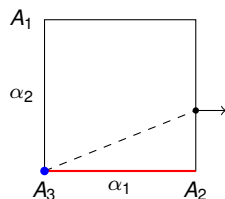
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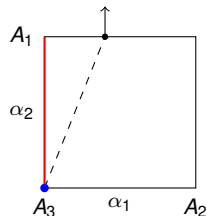
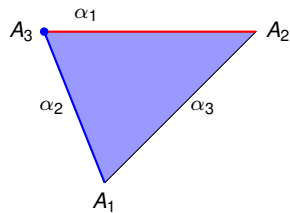
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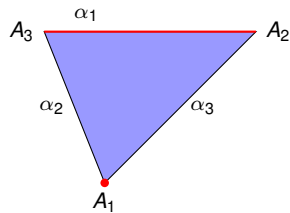
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The reduction process

The self-influence

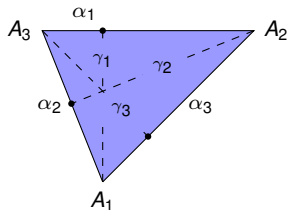
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The reduction process

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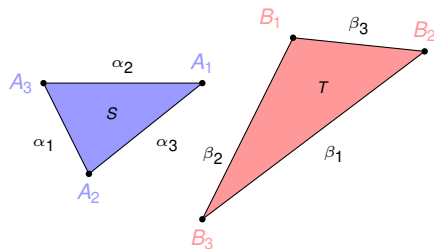


$$I = \frac{2|S|}{3} \sum_{i=1}^3 \gamma_i \left(\arg \sinh \left(\frac{|\alpha_i| - \sigma_i}{\gamma_i} \right) + \arg \sinh \left(\frac{\sigma_i}{\gamma_i} \right) \right)$$

Second case : In the same plane

$$I(S, S) = \frac{2|S|}{3} \sum_{i=1..3} \gamma_i R(A_i, \alpha_i)$$

$$\text{with } R(A_i, \alpha_i) = \int_{\alpha_i} \frac{dx}{\|x - A_i\|}$$

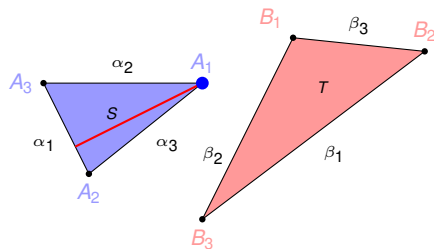


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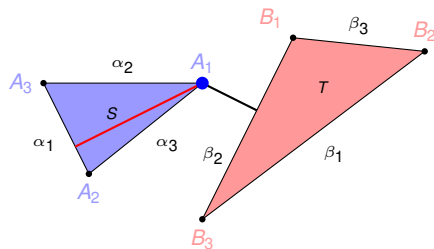


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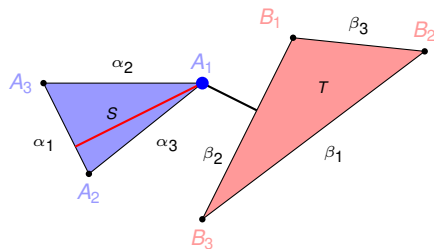


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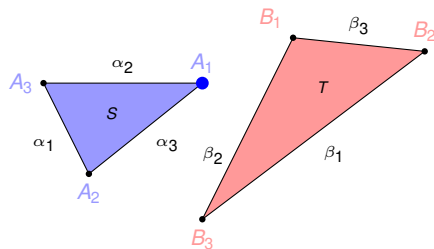


$$I(S, T) = \frac{\gamma_1}{3} \int_{\alpha_1 \times T} \frac{d\gamma_x dy}{\|x - y\|} + \sum_{j=1..3} \frac{\delta_j(A_1)}{3} \int_{S \times \beta_j} \frac{dx d\gamma_y}{\|x - y\|}$$

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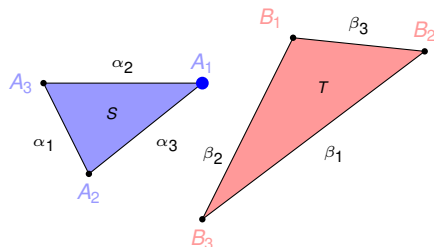


$$I(S, T) = \sum_{i=1..3} \sum_{j=1..3} C_{i,j} R(A_i, \beta_j) + \mathcal{D}_{i,j} R(B_j, \alpha_i)$$

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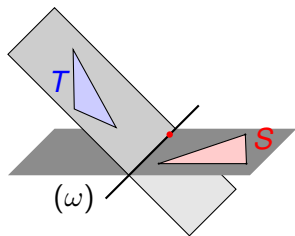
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$$C_{1,2} = \frac{\delta_2(A_1)}{6} \left[\gamma_2(B_3) \left(A_1 - l_{\alpha_2, \beta_2} \left| \frac{A_1 - A_3}{|\alpha_2|} \right| \right) - \gamma_3(B_3) \left(A_1 - l_{\alpha_3, \beta_2} \left| \frac{A_2 - A_1}{|\alpha_3|} \right| \right) \right]$$

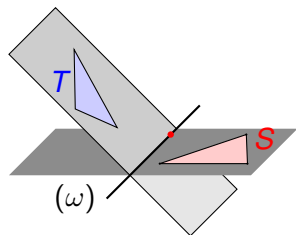
Other cases



We have to choose the origin
in the intersection of the two
planes.

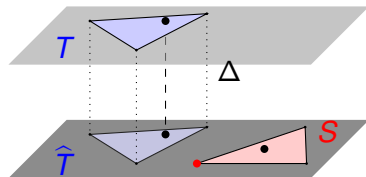
More technical but OK.

Other cases



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More technical but OK.



\hat{T} is the projection of T on the S plane and Δ becomes a parameter.

$$\|x - y\| = (\Delta^2 + \|x - \hat{y}\|^2)^{\frac{1}{2}}$$

! We loose the homogeneity !
can appear with stratified media
or thin plates

What happens

when we lose the homogeneity ?

If the triangles are in parallel planes

$$I = \int_{S \times T} \frac{dx dy}{\|x - y\|} = \int_{S \times \hat{T}} \frac{dx d\hat{y}}{\sqrt{\Delta^2 + \|x - \hat{y}\|^2}}$$

Suppose

- ▶ $\Omega \subset \mathbb{R}^n$
- ▶ $f : (z, \Delta) \in \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ is a homogeneous function of degree r according to z and Δ :

$$f(\lambda z, \lambda \Delta) = \lambda^r f(z, \Delta), \forall \lambda > 0$$

Then

$$\int_{\Omega} f(z, \Delta) dz = \int_{\partial\Omega} (z | \vec{\nu}) \Delta^{r+n} \int_{\Delta}^{+\infty} \frac{f(z, u)}{u^{r+n+1}} du d\gamma_z$$

Proof.

By integration of the Euler's theorem on homogeneous functions and solving an O.D.E. □

What happens

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The second formula

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No reduction of the dimension !

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But if an explicit expression is found for the interior integral

$$\int_{\Omega} f(z, \Delta) dz = \int_{\partial\Omega} (z | \vec{\nu}) F(z, \Delta) d\gamma_z$$

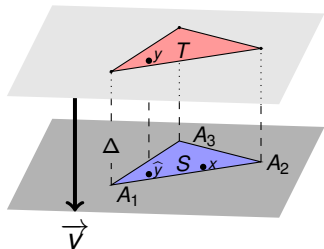
With,

$$F(z, \Delta) = \Delta^{r+n} \int_{\Delta}^{+\infty} \frac{f(z, u)}{u^{r+n+1}} du$$

Superposed Triangles

Double layer potential

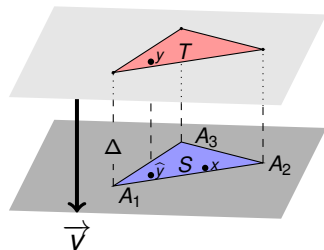
$$\begin{aligned} J &= \int_{S \times T} \frac{x - y}{\|x - y\|^3} dx dy \\ &= \vec{v} \int_{S \times S} \frac{dx d\hat{y}}{(\Delta^2 + \|x - \hat{y}\|^2)^{3/2}} \end{aligned}$$



Superposed Triangles

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After reductions

$$J = 4 \vec{\nu} |S| \sum_{k=1..3} \gamma_i \mathfrak{R}(A_k, \alpha_k, \Delta)$$

Superposed Triangles

$$\Re(A_k, \alpha_k, \Delta) = \left[s \frac{\sqrt{\Delta^2 + \gamma_k^2 + s^2} - \Delta}{2\gamma_k^2(\gamma_k^2 + s^2)} - \frac{s \operatorname{arg} \sinh \left(\frac{\sqrt{\gamma_k^2 + s^2}}{\Delta} \right)}{\gamma_k^2 \sqrt{\gamma_k^2 + s^2}} \right. \\ \left. + \frac{1}{\gamma_k^2} \operatorname{arg} \sinh \left(\frac{s}{\sqrt{\Delta^2 + \gamma_k^2}} \right) + \frac{\gamma_k^2 - \Delta^2}{2\Delta\gamma_k^3} \operatorname{arg} \tan \left(\frac{s}{\gamma_k} \right) - \frac{(\gamma_k^2 - \Delta^2) \pi \operatorname{sgn} s}{4\Delta\gamma_k^3} \right. \\ \left. + \frac{\gamma_k^2 - \Delta^2}{2\Delta\gamma_k^3} \operatorname{Im} \left\{ \operatorname{arg} \tanh \left(\frac{\Delta^2 + \gamma_k^2 + i\gamma_k s}{\Delta \sqrt{\Delta^2 + \gamma_k^2 + s^2}} \right) \right\} \right]_{s^- - \sigma}^{s^+ - \sigma}$$

Δ : the distance between the two planes

γ_k : the distance of A_k to α_k

s^+ , s^- and σ : abscissas

Numerical Results

- Reduction implemented in Matlab (C++ in progress)

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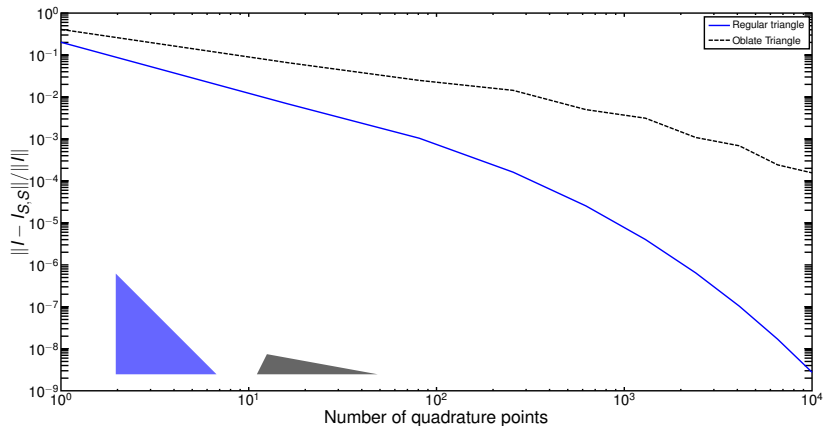
$$I = \int_{S \times S} \frac{dx dy}{\|x - y\|}$$


- We observe our result for the superposed triangles for the double layer potential when $\Delta \rightarrow 0$

$$J = \int_{S \times T} \frac{dx dy}{\|x - y\|^3}$$

Self-influence

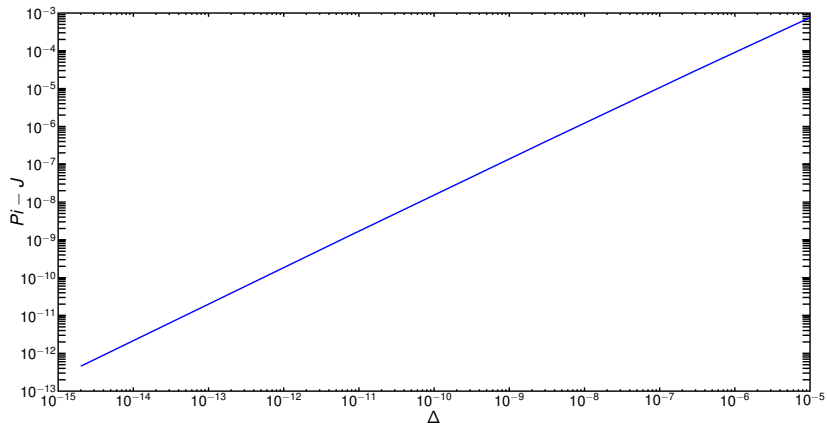
Comparison with Sauter & Schwab's Method



 Sauter, S. and Schwab, C., *Boundary Element Methods*, Springer Series in Computational Mathematics, 2010

Superposed Triangles

The Solid Angle



Conclusion

The method

- Exact formulas for all S and T ; especially for singular and almost singular cases
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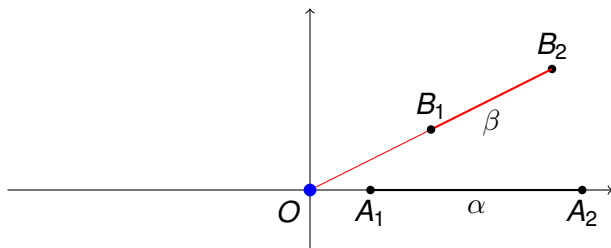
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Work done & Perspectives

- Influence by constant density for single and double layers potentials is treated
- Linear basis function is in progress (self-influence done)
- With linear edge basis functions, Maxwell is partially treated
- Calculation for volumic potential is possible

Thank you for your attention !

The Almost Parallel Case



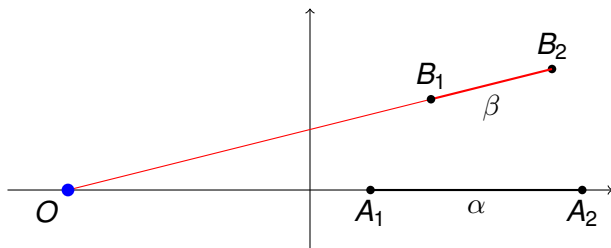
We obtain

$$\begin{aligned} I &= \int_{\alpha \times \beta} \frac{dx dy}{\|x - y\|} \\ &= s_1 R(A_1, \beta) - s_2 R(A_2, \beta) + t_1 R(B_1, \alpha) - t_2 R(B_2, \alpha) \end{aligned}$$

$$\text{with } R(a, \beta) = \int_{\beta} \frac{dx}{\|a - x\|}$$

and $s_1 = \|O - A_1\|$, $t_1 = \|O - B_1\| \dots$

The Almost Parallel Case



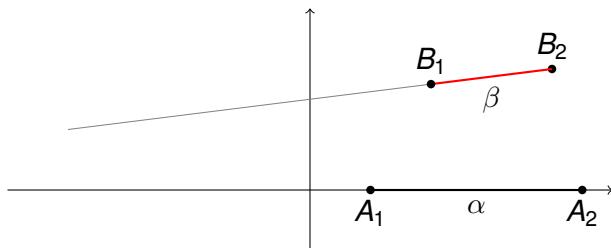
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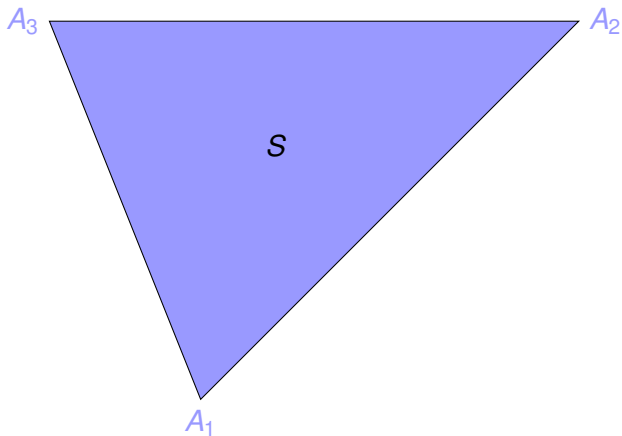
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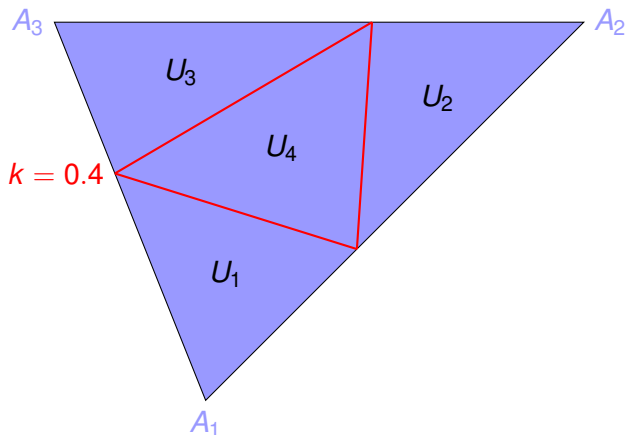
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Almost parallel case



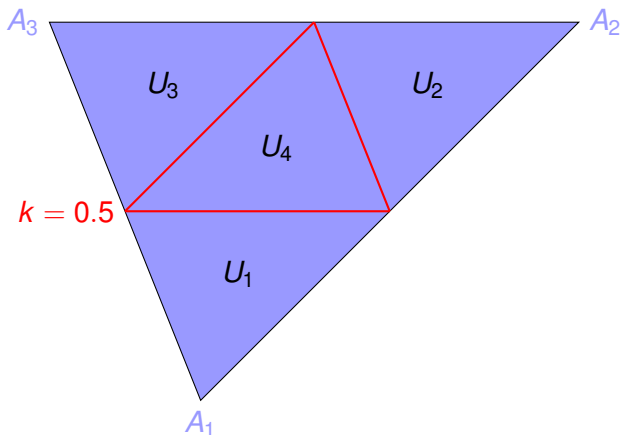
$$I = \int_{S \times S} \frac{dxdy}{\|x - y\|}$$

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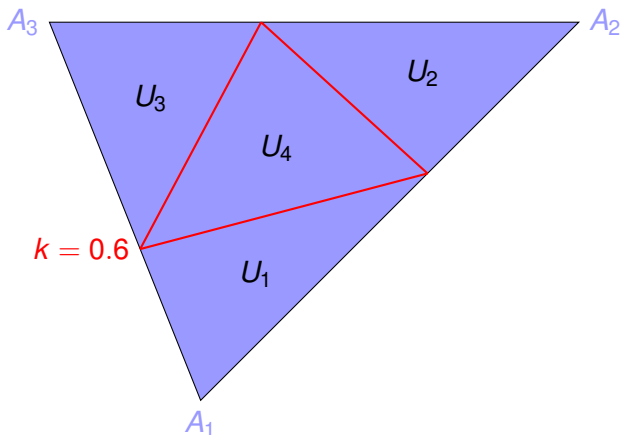
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